

Improving the security of quantum exam against cheating

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Abstract

The security of quantum exam [Phys. Lett. A 350 (2006) 174] is analyzed and it is found that this protocol is secure for any eavesdropper except for the “students” who take part in the exam. Specifically, any student can steal other examinees’ solutions and then cheat in the exam. Furthermore, a possible improvement of this protocol is presented.

Key words: quantum cryptography, cryptanalysis, entanglement

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Cryptography is the approach to assure the secrecy of the data which is stored or communicated in public environment. From its beginning the research of cryptography has been progressed along two directions in parallel. One direction deals with the design of various schemes to maintain privacy. The other is focused on analyzing the security of existing protocols, trying to find the flaws in cryptosystems and improve them. Both directions are necessary to the development of cryptography. It is also the case in quantum cryptography [1,2,3], where the work of both scheme designing (e.g. [4] and references therein) and security analyzing (e.g. [5,6,7,8]) is continually proposed.

In a recent paper [9] a novel protocol called quantum exam was proposed. In this protocol a teacher Alice wants to organize an exam with her remotely separated students Bob 1, Bob 2, ... and Bob N . As in a classical exam, all the problems and Bobs’ solutions should not be leaked out and, more importantly, any Bob cannot obtain other examinees’ solutions. However, we

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find that the later confidentiality constraint is not perfectly satisfied. That is, a dishonest Bob can cheating in the exam. In this Letter we demonstrate this hidden trouble and then present a possible improvement of the quantum exam protocol.

Let us introduce the quantum exam first. In fact there are two similar quantum exam protocols presented in Ref.[9]. We will take the first one (i.e., the so-called absolutely secure protocol) as our example. For simplicity we use the same notations as that in Ref.[9]. The whole protocol is a little complicated and here we only describe briefly the related part, that is, the solution-collecting part (including the entanglement-sharing process). In this stage Alice generates a large enough number of ordered nonidentical states

$$|\Phi_p\rangle_{a_p 1_p \dots N_p} = \frac{1}{\sqrt{2}}(|0s_{1_p}s_{2_p}\dots s_{N_p}\rangle_{a_p 1_p \dots N_p} + |1\bar{s}_{1_p}\bar{s}_{2_p}\dots\bar{s}_{N_p}\rangle_{a_p 1_p \dots N_p}), \quad (1)$$

where $s_{n_p} = 0$ or 1 , $\forall 1 \leq n \leq N$, and $\bar{s}_{n_p} = s_{n_p} \oplus 1$ (\oplus denotes an addition mod 2). Note that the value of s_{n_p} is known only to Alice. For each $|\Phi_p\rangle$ Alice stores qubit a_p and sends qubits $1_p, 2_p, \dots, N_p$ to Bob 1, Bob 2, ..., Bob N , respectively. Afterwards, Alice selects a subset of the entangled states $\{|\Phi_l\rangle\}$ to detect eavesdropping. More concretely, for each $|\Phi_l\rangle$, Alice measures the qubit a_l randomly in the basis B_z or B_x and informs every Bob to perform the same measurement on his corresponding qubit. Then they check the security of the entanglement distribution process by verifying

$$j_{a_l}^z = \delta_{0,s_{n_l}} j_{n_l}^z + \delta_{1,s_{n_l}} (j_{n_l}^z \oplus 1) \quad (2)$$

for every $n = 1, 2, \dots, N$ (when B_z was used) or

$$j_{a_l}^x = \prod_{n=1}^N j_{n_l}^x \quad (3)$$

(when B_x was used), where j represents the measurement result, $j_{a_l}^z(j_{n_l}^z) = \{0, 1\}$ corresponding to obtaining $\{|0\rangle, |1\rangle\}$ and $j_{a_l}^x(j_{n_l}^x) = \{+1, -1\}$ corresponding to obtaining $\{|+\rangle, |-\rangle\}$. If there is no eavesdropping detected, the shared entanglement can be used for solution-collecting some time later. When needed, Alice and Bobs measure the remaining ordered $|\Phi_p\rangle$ -states $\{|\Phi_m\rangle_{a_m 1_m \dots N_m}\}$ in basis B_z and record the outcomes as the secure keys. Let $\{j_{a_m}^z\}$ and $\{j_{n_m}^z\}$ denote the keys belonging to Alice and Bob n , respectively. Every Bob uses his key as a one-time-pad to encrypt his solution and sends it to Alice. With the knowledge of $j_{a_m}^z$ and s_{n_m} Alice can obtain each Bob's key [see Eq.(2)]. Consequently, at the end of the exam Alice will correctly decrypt Bobs' messages and obtain every Bob's solution.

It can be seen that the solution-collecting process comprises mainly a multipartite quantum key distribution (MQKD) scheme. Because the one-time-pad is perfectly secure here, the security of the whole process lies on that of the key distribution. As we know, the state $|\Phi_p\rangle_{a_p 1_p \dots N_p}$ has a property of positive parity, i.e., $j_{a_p}^x \prod_{n=1}^N j_{n_p}^x = +1$. This wonderful property is subtly employed to detect eavesdropping in the quantum exam protocol [see Eq.(3)]. As a result, the two constraints Eqs.(2) and (3) can make the exam secure against various kinds of attacks [9]. However, we take notice of another property of $|\Phi_p\rangle_{a_p 1_p \dots N_p}$, that is, one can entangle an ancilla $|0\rangle$ into the multipartite entangled state by a controlled-NOT (CNOT) operation and then disentangle it out from the obtained state by another CNOT operation. The control qubits of the two CNOT operations can be any two qubits in $|\Phi_p\rangle_{a_p 1_p \dots N_p}$ and the target is the ancilla. For example, for a certain p , the multipartite entangled state and the ancilla compose a composite system

$$|\Gamma\rangle^1 = |\Phi\rangle_{a1\dots N}|0\rangle_g = \frac{1}{\sqrt{2}}(|0s_1s_2\dots s_N\rangle_{a1\dots N}|0\rangle_g + |1\bar{s}_1\bar{s}_2\dots\bar{s}_N\rangle_{a1\dots N}|0\rangle_g), \quad (4)$$

where the subscript g represents the ancilla. If one performs a CNOT operation C_{kg} (the first subscript k denotes the control qubit and the second one g denotes the target qubit) on the qubit k ($1 \leq k \leq N$) and the ancilla, the state of the system changes into

$$|\Gamma\rangle^2 = \frac{1}{\sqrt{2}}(|0s_1s_2\dots s_N\rangle_{a1\dots N}|s_k\rangle_g + |1\bar{s}_1\bar{s}_2\dots\bar{s}_N\rangle_{a1\dots N}|\bar{s}_k\rangle_g). \quad (5)$$

Now if one performs another CNOT operation C_{rg} on the qubit r ($1 \leq r \leq N$) and the ancilla, he (she) will obtain

$$\begin{aligned} |\Gamma\rangle^3 &= \frac{1}{\sqrt{2}}(|0s_1s_2\dots s_N\rangle_{a1\dots N}|s_k \oplus s_r\rangle_g + |1\bar{s}_1\bar{s}_2\dots\bar{s}_N\rangle_{a1\dots N}|\bar{s}_k \oplus \bar{s}_r\rangle_g) \\ &= \frac{1}{\sqrt{2}}(|0s_1s_2\dots s_N\rangle_{a1\dots N}|s_k \oplus s_r\rangle_g + |1\bar{s}_1\bar{s}_2\dots\bar{s}_N\rangle_{a1\dots N}|s_k \oplus s_r\rangle_g) \\ &= |\Phi\rangle_{a1\dots N}|s_k \oplus s_r\rangle_g. \end{aligned} \quad (6)$$

It can be seen that the ancilla is disentangled out from the multipartite entangled state and, more importantly, the original state $|\Phi\rangle_{a1\dots N}$ is left alone. As a result, if an eavesdropper Eve utilizes the above operations to eavesdrop, she will introduce no errors. Furthermore, when Eve measures the ancilla in basis B_z she will obtain $s_k \oplus s_r$ definitely. Since the value $s_k \oplus s_r$ implies, as described as following, the correlation of the measurement results of qubits k and r , we call the state $|\Phi\rangle_{a1\dots N}$ “correlation elicitable”. It can be shown that this property gives a dishonest Bob the chance to cheat in the exam. Without loss of generality, suppose the dishonest student is Bob r and he wants to steal

Bob k 's solution (maybe Bob k is an outstanding student), he can adopt the following strategy to achieve his goal.

- (i) For each p , Bob r prepares an ancilla $|0\rangle$ and performs two CNOT operations $C_{k_pg_p}$ and $C_{r_pg_p}$ as described above when Alice distributes the multipartite entangled states $\{|\Phi_p\rangle_{a_p1_p\dots N_p}\}$.
- (ii) Bob r measures each ancilla in basis B_z and obtains $s_{k_p} \oplus s_{r_p}$ with certainty.
- (iii) Cooperating with Alice, Bob r executes the legal process to detect eavesdropping and get key bits. After the actions (i) and (ii), as analyzed above, all the carrier states $\{|\Phi_p\rangle_{a_p1_p\dots N_p}\}$ remain unchanged and no disturbance is introduced. Therefore, Alice cannot detect the eavesdropping and Bob r will correctly obtain the intended key bits $\{j_{r_m}^z\}$.
- (iv) Bob r gains Bob k 's key bits $\{j_{k_m}^z\}$ by simple calculation. More specifically, Bob r deletes the data corresponding to the check states $\{|\Phi_l\rangle\}$ from the bits $\{s_{k_p} \oplus s_{r_p}\}$, and obtains the remaining ordered bits $\{s_{k_m} \oplus s_{r_m}\}$, which correspond to the carrier states $\{|\Phi_m\rangle_{a_m1_m\dots N_m}\}$ and the key bits $\{j_{r_m}^z\}$. It should be emphasized that, for a certain m , the measurement outcomes of the ancilla $s_{k_m} \oplus s_{r_m}$ implies the relation between two key bits $j_{k_m}^z$ and $j_{r_m}^z$, that is, $j_{k_m}^z \oplus j_{r_m}^z = s_{k_m} \oplus s_{r_m}$. [From Eq.(1) we can see that either $j_{k_m}^z = s_{k_m}, j_{r_m}^z = s_{r_m}$ or $j_{k_m}^z = \bar{s}_{k_m}, j_{r_m}^z = \bar{s}_{r_m}$ holds.] Therefore, with the knowledge of $\{s_{k_m} \oplus s_{r_m}\}$ and $\{j_{r_m}^z\}$, Bob r can easily get the key bits $\{j_{k_m}^z\}$ of Bob k by calculating $j_{k_m}^z = s_{k_m} \oplus s_{r_m} \oplus j_{r_m}^z$ for each m .
- (v) Bob r cheats when Alice collects the solutions. Obviously, with the help of $\{j_{k_m}^z\}$, Bob r can decrypt the message sent from Bob k to Alice and copy Bob k 's solution at will.

By this strategy, a dishonest student can steal any other examinees' solutions. Moreover, the eavesdropping is not difficult to realize because it needs only facilities similar to that of the legal parties. One may argue that, in the above example, if Bob r is far away from the quantum channel between Alice and Bob k he cannot continually perform the two CNOT operations in a certain time. In fact there is no need to worry about it. Bob r does not need to take a round trip between his and Bob k 's quantum channels. He can ask his friend, say Charlie, who stands in Bob k 's channel, to perform the first CNOT operation $C_{k_pg_p}$ and then send the ancilla to him.

There is a fact which should be pointed out. That is, the one who will legally take part in the protocol is prone to be omitted when we analyze various attack strategies. In fact, in most MQKD protocols (e.g. quantum secret sharing, see [10] and references therein), a participant generally has more power to attack than an outside eavesdropper because the participant can take advantage of the right to access the carrier state partly and participate in the process of

eavesdropping detection. We call this kind of attack “participant attack”. In the quantum exam protocols, as we can see, the eavesdropping result $\{s_{k_m} \oplus s_{r_m}\}$ does not seem to have much meaning for an outside eavesdropper, but it is very useful for a participant Bob to eavesdrop further. Therefore, as implied in Refs.[11,12,13], the main goal for the security of an MQKD should be focused on preventing the dishonest participant from eavesdropping the information.

Now we discuss how to improve the quantum exam protocol to prevent this kind of participant attack. To retain the features of the original quantum exam protocol, our aim is to modify it as little as possible. Since the fundamental reason of this threat is the speciality of $|\Phi\rangle_{a1\dots N}$, i.e., “correlation elicitable”, Alice can insert some different check qubits to detect the above attack. For example, before Alice sends the sequences to Bobs, she inserts a certain number of single qubits into each sequence in random positions. All these single qubits are randomly in one of the states $\{|+\rangle, |-\rangle\}$ [14]. Note that the positions of the single qubits in these sequences are different from each other. After all Bobs received their respective sequences, Alice tells each Bob the positions of these check qubits and lets him measure them in the basis B_x . Then Alice and Bob check the identity of these qubits. If the error rate is low enough, they proceed with other steps in the original protocol to finish the quantum exam. Because, for the dishonest Bob, both the single qubits and the qubits from $|\Phi\rangle_{a1\dots N}$ are in maximally mixed state $\rho = \frac{1}{2}(|0\rangle\langle 0| + |1\rangle\langle 1|)$, he cannot distinguish the check qubits from others. Therefore, when the dishonest Bob wants to cheat using above strategy, he would introduce errors with probability $\frac{1}{2}$ once he performs a CNOT operation on a certain check qubit and his ancilla. As a result, the improved protocol can stand against the above participant attack. Furthermore, the main frame of the original protocol is retained and it follows that the security against other kinds of attacks (such as measure-resend attack, disturbance attack, entangle-measure attack, etc. [9]) still holds.

In conclusion, we show that a dishonest student can cheat in the quantum exam [9] and give a possible improvement by inserting some additional check qubits in each sequence. We emphasize that the participant attack should not be overlooked when we discuss the security of a MQKD scheme, which generally possesses more power in eavesdropping than the attack from outside.

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- [13] F. G. Deng, X. H. Li, H. Y. Zhou, *et al.*, Phys. Rev. A 72 (2005) 044302.
- [14] Here the role of the states $\{|+\rangle, |-\rangle\}$ is just to prevent the presented attack. To acquire more security the original strategy to detect eavesdropping is still needed. We can also use four states $\{|0\rangle, |1\rangle, |+\rangle, |-\rangle\}$ to prepare the inserted single qubits, which can totally ensure the security of these sequences (similar with that of BB84 protocol [1]). We do not choose the latter choice because we try to retain the features of the original quantum exam protocol, including its strategy to detect eavesdropping.